Crack propagation in heterogeneous materials with several defects

A. Piccolroaz¹ G. Mishuris¹ A. Movchan² N. Movchan²

1 Institute of Mathematics and Physics Aberystwyth University and Wales Institute of Mathematical and Computational Sciences

> ²Department of Mathematical Sciences University of Liverpool

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The problem and motivation

Crack advancing in a bimaterial plane with finite array of defects

- What is the effect of the defects on the propagation of the crack? Amplification/shielding effects?
- Can we arrange the defects in such a way to stop the propagation of the crack?
- Vice versa, can we arrange the defects in such a way to make the crack propagate up to the end of the array?

New challenges:

- Interfacial crack
- Weight function for an interfacial crack
- Singular perturbation procedure related to small defects
- Singular perturbation procedure related to small advance of the crack

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Type of defects

- **•** Small elastic inclusions
- **•** Small voids
- Small rigid inclusions
- **•** Microcracks
- Rigid line inclusions

Line defect with imperfect bonding:

• Line defect with soft bonding (stiffness κ): $\llbracket \sigma \rrbracket(s) = 0, \quad \sigma(s) = \kappa \llbracket u \rrbracket(s)$

 $\int \kappa = 0 \implies \sigma(s) = 0$ microcrack $\kappa = \infty \quad \Rightarrow \quad [\![u]\!](s) = 0 \quad$ perfect bonding (no defect)

Stiff line defect (stiffness κ): $\left[\![u]\!](s) = 0, \quad \left[\![\sigma]\!](s) + \kappa \left.\frac{\partial^2 u}{\partial s^2}\right.\right|_{\gamma^{\epsilon}} = 0$

$$
\begin{cases}\n\kappa = 0 & \Rightarrow \quad \lbrack \sigma \rbrack(s) = 0 \quad \text{no defect} \\
\kappa = \infty & \Rightarrow \quad \frac{\partial^2 u}{\partial s^2}\Big|_{\gamma \infty} = 0 \quad \text{rigid line inclusion}\n\end{cases}
$$

Problem formulation

Bimaterial plane with a dominant crack along the interface and small defects:

Singular perturbations:

- **1** Elastic inclusion
- **2** Microcrack
- ³ Rigid line inclusion

Singular perturbation:

1 Crack advance

Small parameter ε : diameter of defect $2\varepsilon l, \varepsilon \ll 1$

Assumptions:

- **1** Mode III deformation
- ² Loading on the crack surfaces
- ³ Perfect interfaces (continuity of displacements and tractions)
- The composite is dilute (neglect interactions between small defects)
- ⁵ Stable quasi-static propagation (neglect inertia term[s\)](#page-2-0) \Box

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$$
u(\mathbf{x},\varepsilon) = u^{(0)}(\mathbf{x}) + \left[\varepsilon \sum_{j=1}^3 W_j(\boldsymbol{\xi}_j)\right] + \varepsilon^2 \sum_{j=1}^3 u^{(j)}(\mathbf{x}) + \varepsilon^2 v(\mathbf{x},\phi) + o(\varepsilon^2), \quad \varepsilon \to 0
$$

1 Solution of the unperturbed problem ($\varepsilon = 0$) ² Boundary layers concentrated near the defects

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³ Additional terms to adjust the BC and IC disturbed by the boundary layers

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- Perturbation associated with the crack advance $\varepsilon^2\phi$

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1 Solution of the unperturbed problem ($\varepsilon = 0$) Boundary layers concentrated near the defects

- ³ Additional terms to adjust the BC and IC disturbed by the boundary layers
- Perturbation associated with the crack advance $\varepsilon^2\phi$
- Using the linearity of the problem, we analyse the perturbation of each defect separately (superposition principle).
- \bullet The method can be extended to a finite number of defects, provided that the distance between defects remains finite (composite is dilute).

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The first challenge: Weight functions for an interfacial crack

"Weight functions are fundamental solutions for the problem of a cracked body"

Linear elasticity:

Linear fracture mechanics:

Interfacial crack:

Green's function for elastic body

Point forces **SIFs** *a*

Weight functions for cracked homogeneous body

 $u_i(x, y) = G_{ii}(x, y)e_i$

$$
K = \sqrt{\frac{2}{\pi}} (1 + i) a^{-1/2}
$$

Weight functions for interfacial crack

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The weight functions for an interfacial crack

Mode III symmetric and skew-symmetric weight functions:

$$
\llbracket U_3 \rrbracket(x_1) = \begin{cases} \frac{1 - i}{\sqrt{2\pi}} x_1^{-1/2}, & \text{for } x_1 > 0, \\ 0, & \text{for } x_1 < 0, \end{cases} \qquad \langle U_3 \rangle = \frac{\eta}{2} \llbracket U_3 \rrbracket,
$$

• Integral formula for the computation of the Mode III SIF:

$$
K_{\rm III} = -(1+i)\lim_{x_1'\to 0^+} \int_{-\infty}^0 \{\underbrace{\llbracket U_3 \rrbracket (x_1'-x_1) \langle p_3 \rangle (x_1)}_{\text{symmetric part}} + \underbrace{\langle U_3 \rangle (x_1'-x_1) \llbracket p_3 \rrbracket (x_1)}_{\text{skew-symmetric part}} \} dx_1
$$

$$
\eta = \frac{\mu - \mu_+}{\mu_- + \mu_+}
$$
 bimaterial parameter

 $= \Omega$

 $\mathcal{A} \equiv \mathcal{B} \times \mathcal{A} \equiv \mathcal{B}$

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The second challenge: singular perturbation for a small defect

The defect is replaced by "effective" tractions

Dipole field:

$$
\varepsilon^2 w(\mathbf{x}) = -\frac{\varepsilon^2}{2\pi} \left[\nabla u^{(0)} \Big| \mathbf{y} \right] \cdot \left[\mathbf{\mathcal{M}} \frac{\mathbf{x} - \mathbf{Y}}{|\mathbf{x} - \mathbf{Y}|^2} \right] + o(\varepsilon^2), \quad \varepsilon \to 0
$$

 M is the dipole matrix

Perturbation of the SIF:

$$
\left(\Delta K_{\parallel\parallel} = -\sqrt{\frac{2}{\pi}} \frac{\mu_{+} \mu_{-}}{\mu_{+} + \mu_{-}} \nabla u^{(0)}\Big|_{\mathbf{Y}} \cdot \mathbf{M} \mathbf{c}\right)
$$

$$
\mathbf{c} = \frac{1}{2d^{3/2}} \left[-\sin \frac{3\varphi}{2}, \cos \frac{3\varphi}{2} \right]
$$

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Example: Elliptic elastic inclusion

Dipole field:

$$
\varepsilon^2 w(\mathbf{x}) = -\frac{\varepsilon^2}{2\pi} \left[\nabla u^{(0)} \Big| \mathbf{y} \right] \cdot \left[\mathbf{\mathcal{M}} \frac{\mathbf{x} - \mathbf{Y}}{|\mathbf{x} - \mathbf{Y}|^2} \right] + o(\varepsilon^2), \quad \varepsilon \to 0
$$

Dipole matrix

$$
\mathcal{M} = -\frac{\pi}{2}ab(1+e)(\mu_{\star}-1)\left[\begin{array}{cc} \frac{1+\cos 2\alpha}{e+\mu_{\star}} + \frac{1-\cos 2\alpha}{1+e\mu_{\star}} & -\frac{(1-e)(\mu_{\star}-1)\sin 2\alpha}{(e+\mu_{\star})(1+e\mu_{\star})} \\ -\frac{(1-e)(\mu_{\star}-1)\sin 2\alpha}{(e+\mu_{\star})(1+e\mu_{\star})} & \frac{1-\cos 2\alpha_{1}}{e+\mu_{\star}} + \frac{1+\cos 2\alpha}{1+e\mu_{\star}} \end{array}\right]
$$

$$
\mu_\star = \mu/\mu_i \qquad e = b/a
$$

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Dipole matrix for different types of defects

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Application: Shielding and amplification effects

$$
\Delta K_{\text{III}} = -\sqrt{\frac{2}{\pi}} \frac{\mu_{+} \mu_{-}}{\mu_{+} + \mu_{-}} \nabla u^{(0)} \bigg|_{\mathbf{Y}} \cdot \mathbf{M} \mathbf{c}
$$

Definition:

∆*K*III < 0: *shielding effect* "the defect is *preventing* the propagation" ∆*K*III > 0: *amplification effect* "the defect is *promoting* the propagation" $\Delta K_{\text{III}} = 0$: *neutral* "the defect has no effect"

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Example: shielding/amplification diagrams for macro-microcrack interaction

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The third challenge: Singular perturbation for crack advance

The Betti identity:

$$
\int_{-\infty}^{\infty} \left\{ \llbracket U \rrbracket (x_1' - x_1) \langle \sigma \rangle (x_1) + \langle U \rangle (x_1' - x_1) \llbracket \sigma \rrbracket (x_1) - \langle \Sigma \rangle (x_1' - x_1) \llbracket u \rrbracket (x_1) \right\} dx_1 = 0
$$

 u, σ physical solution *U*, Σ weight functions

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Perturbation of the SIF:

$$
\widehat{\left(\varepsilon^2 \Delta K_{\text{III}}^{\phi} = \frac{\varepsilon^2 \phi}{2} A_{\text{III}}^{(0)}\right)}
$$

$$
A_{\parallel\parallel}^{(0)} = \sqrt{\frac{2}{\pi}} \int_{-\infty}^{0} \left(\langle p \rangle(x_1) + \frac{\eta}{2} [p] [(x_1) \right) (-x_1)^{-3/2} dx_1
$$

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Analysis of a stable quasi-static propagation

The stress intensity factor is expanded as follows:

$$
K_{\parallel\parallel} = K_{\parallel\parallel}^{(0)} + \varepsilon^2 \left(\Delta K_{\parallel\parallel}^{\phi} + \sum_{j=1}^3 \Delta K_{\parallel\parallel}^{(j)} \right) + o(\varepsilon^2), \quad \varepsilon \to 0
$$

 $\Delta K_{\parallel \parallel}^{\phi} = \frac{\varepsilon^2 \phi}{2}$ $\frac{\varphi}{2}A_{\textrm{III}}^{(0)}$: perturbation produced by the
elongation of the crack along the interface

 $\sum_{j=1}^{3} \Delta K_{\text{III}}^{(j)}$: perturbation produced by the defects

We assume that the crack propagation is stable and quasi-static: $G = G_c$

$$
G = \frac{1}{4} \left(\frac{1}{\mu_{+}} + \frac{1}{\mu_{-}} \right) K_{III}^{2} \Rightarrow \Delta K_{III}^{\phi} + \sum_{j=1}^{3} \Delta K_{III}^{(j)} = 0 \Rightarrow \varphi = \frac{2}{A_{III}^{(0)}} \sum_{j=1}^{3} \Delta K_{III}^{(j)}
$$

$$
\Delta K_{\parallel \parallel}^{(j)} = -\sqrt{\frac{2}{\pi}} \frac{\mu_{+} \mu_{-}}{\mu_{+} + \mu_{-}} \nabla u^{(0)} \Big|_{\mathbf{Y}_{j}} \cdot \mathbf{\mathcal{M}}_{j} \mathbf{c}_{j}, \qquad \mathbf{c}_{j} = \frac{1}{2d_{j}^{3/2}} \left[-\sin \frac{3\varphi_{j}}{2}, \cos \frac{3\varphi_{j}}{2} \right]
$$

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Application: Crack propagation and arrest

2εl₃ Given a configuration of defects and position of the crack tip, the incremental crack advance ϕ is given by:

$$
\phi = \frac{2}{A_{\text{III}}^{(0)}} \sum_{j=1}^{3} \Delta K_{\text{III}}^{(j)}
$$

It is possible to update the configuration with the new position of the crack tip and recompute the incremental crack advance in the new configuration, following an iterative procedure:

- \bullet the crack "accelerates" when the increment ϕ is increasing
- \bullet the crack "decelerates" when the increment ϕ is decreasing
- **the crack "arrests" when a neutral configuration is reached (** $\phi = 0$ **)**

The total crack elongation is computed as:

$$
x(N) = \sum_{i=0}^{N} \varepsilon^2 \phi_i
$$

where *N* is the number of iterations.

Crack propagation in a finite array of defects

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Conclusions and possible extensions

This asymptotic method can be extended to the case of an infinite array of defects

-4 *w* -2 *w* 0 2 *w* 4 *w* -2 *h* -*h* 0 *h* 2 *h x y* − *j = −N , ... ,* −1 *j =* 1, ... , *N* + *j =* −2 0 1 2 3 4 5 −1 *h w* −3

$$
\phi = \frac{al^2 \cos 2\alpha}{2h^2} \left\{ -1 + \sum_{j=1}^{N^+} \frac{\left(1 - \frac{h^2}{j^2 w^2}\right) \frac{h^2}{j^2 w^2}}{\left(1 + \frac{h^2}{j^2 w^2}\right)^2} + \sum_{j=1}^{N^-} \frac{\left(1 - \frac{h^2}{j^2 w^2}\right) \frac{h^2}{j^2 w^2}}{\left(1 + \frac{h^2}{j^2 w^2}\right)^2} \right\}
$$

Take the limit as $N^+ \to \infty$:

$$
\phi = \frac{al^2 \cos 2\alpha}{2h^2} \left\{ -\frac{1}{2} - \frac{\left(\frac{\pi h}{w}\right)^2}{2 \sinh^2\left(\frac{\pi h}{w}\right)} + \sum_{j=1}^{N^-} \frac{\left(1 - \frac{h^2}{j^2 w^2}\right) \frac{h^2}{j^2 w^2}}{\left(1 + \frac{h^2}{j^2 w^2}\right)^2} \right\}
$$

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References

N.J. Hardiman. Elliptic elastic inclusion in an infinite elastic plate. *QJMAM* **7**, 226–230, 1954.

H.F. Bueckner. A novel principle for the computation of stress intensity factors. *ZAMM* **46**, 529–545, 1970.

A. Piccolroaz, G. Mishuris & A.B. Movchan. Symmetric and skew-symmetric weight functions in 2D perturbation models for semi-infinite interfacial cracks. *JMPS* **57**, 1657–1682, 2009.

G. Mishuris, A. Movchan, N. Movchan & A. Piccolroaz. Interaction of an interfacial crack with linear small defects under out-of-plane shear loading. *CMS*, in press.

A. Piccolroaz, G. Mishuris, A. Movchan & N. Movchan. Perturbation analysis of Mode III interfacial crack advancing in a dilute heterogeneous material. Submitted.

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